

Acoustic Pulses

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Acoustic Wave Equation (compressible gas)

$$\left(\vec{\partial}^2 - \frac{1}{c^2}\partial_t^2\right)\phi(\vec{r}, t) = 0 \quad (1)$$

Here ϕ is the velocity potential, i.e., $\vec{\partial}\phi(\vec{r}, t) = \vec{v}(\vec{r}, t)$ where $\vec{v}(\vec{r}, t)$ is the velocity of the gas atoms/molecules at position \vec{r} and time t . c is the speed of sound.

Consider the 1+1 dimensional problem where an infinite inflexible movable "wall" is placed parallel to the yz plane with it's position along the x axis given by $X(t)$. Let the gas fill the semi-infinite half space to the right of the wall, i.e., $x > X(t)$. When the wall speed $\partial_t X(t)$ is much less than the speed of sound the atoms/molecules just to the right of the wall must move with the speed of the wall, i.e. $v(X(t), t) = \partial_t X(t)$. This is only one boundary condition, nominally we need two. But since the wave equation can be factorized as $(\partial_x - \partial_t/c)(\partial_x + \partial_t/c)\phi = 0$ and we can solve either $(\partial_x - \partial_t/c)\phi = 0$ or $(\partial_x + \partial_t/c)\phi = 0$ to get waves/pulses moving in the $-x$ or the $+x$ direction, respectively we only need one boundary condition which is $(\partial_x\phi)(X(t), t) = v(X(t), t) = \partial_t X(t)$.

For $\partial_t X(t) \ll c$ we have the following approximate solution

$$\phi(x, t) = -cX\left(t - \frac{x - X(t)}{c}\right) \quad (2)$$

First show that this satisfies the wave equation up to terms of order $\partial_t X/c$. Let $\xi(x, t) = t - (x - X(t))/c$, then

$$\begin{aligned} \left(\partial_x + \frac{1}{c}\partial_t\right)(-cX(\xi(x, t))) &= -c(\partial_\xi X)(\partial_x \xi) + \frac{1}{c}(-c)(\partial_\xi X)(\partial_t \xi) \\ &= -c(\partial_\xi X)\left(-\frac{1}{c}\right) - (\partial_\xi X)\left(1 + \frac{\partial_t X(t)}{c}\right) \end{aligned}$$

But given $\partial_t X(t)/c \ll 1$ we have

$$\left(\partial_x + \frac{1}{c}\partial_t\right)(-cX(\xi(x, t))) = \left(\partial_x + \frac{1}{c}\partial_t\right)\left(-cX\left(t - \frac{x - X(t)}{c}\right)\right) \cong 0 \quad (3)$$

Show that it satisfies the boundary conditions.

$$\partial_x \phi(x, t)|_{x=X(t)} = \partial_\xi X(\xi(x, t))|_{x=X(t)} = (\partial_\xi X)(t) = \partial_t X(t) \quad (4)$$

The density variation $\delta\rho(x, t)$ about a constant background density, ρ_0 , is related to the speed of the molecules via

$$\frac{v(x, t)}{c} = \frac{\delta\rho(x, t)}{\rho_0} \quad (5)$$

and so

$$\frac{\delta\rho(x, t)}{\rho_0} = \frac{\partial_x \phi(x, t)}{c} = -\partial_x X\left(t - \frac{x - X(t)}{c}\right) \quad (6)$$

For the wall moving at a constant speed, $X(t) = Vt$, we have

$$\frac{\delta\rho(x, t)}{\rho_0} = -\partial_x\left(V \times \left(t - \frac{x - Vt}{c}\right)\right) = \frac{V}{c} \quad (7)$$

All molecules are moving at the same speed, $\partial_x v(x, t) = 0$. In order to generate acoustic waves or pulses the wall must change speed, i.e., accelerate.

Consider the case where the wall starts with zero speed and accelerates up to a finite speed. For a step function jerk profile, the acceleration, velocity and position look as follows. NOTE: jerk is the derivative of acceleration with respect to time and so a step function jerk induces a linear time dependence of the acceleration.



